

Supplementary Material to the article

“Controlling the dynamic modes of a nanoscale superconducting interferometer”

In this supplementary material we derive the solution of the second order inhomogeneous differential equation:

$$\ddot{\varphi} + \alpha\dot{\varphi} + (1 + \varepsilon \text{Cos}(2\omega_0\tau + \Delta))\omega_0^2\varphi = i_s \text{Cos}(\omega_0\tau). \quad (\text{S1})$$

The general solution of the equation (S1) is the sum of the general solution of a homogeneous equation and the particular solution of an inhomogeneous equation:

$$\varphi(\tau) = \varphi_{gshe}(\tau) + \varphi_{psie}(\tau). \quad (\text{S2})$$

We assume the general solution of the homogeneous equation as $\varphi_{gshe}(\tau) = \mu e^{k\tau} \text{Cos}(\omega_0\tau) + \nu e^{k\tau} \text{Sin}(\omega_0\tau)$. After the substituting of $\varphi_{gshe}(\tau)$ into the equation (S1) the following coupled linear equations for the coefficients μ and ν were obtained:

$$\begin{aligned} \nu \left(2\omega_0 k + \alpha\omega_0 - \frac{\omega_0^2 \varepsilon}{2} \text{Sin}(\Delta) \right) \\ + \mu \left(k\alpha + \frac{\omega_0^2 \varepsilon}{2} \text{Cos}(\Delta) \right) \\ = 0, \end{aligned} \quad (\text{S3a})$$

$$\begin{aligned} \mu \left(-2\omega_0 k - \alpha\omega_0 - \frac{\omega_0^2 \varepsilon}{2} \text{Sin}(\Delta) \right) \\ + \nu \left(k\alpha - \frac{\omega_0^2 \varepsilon}{2} \text{Cos}(\Delta) \right) \\ = 0. \end{aligned} \quad (\text{S3b})$$

The compatibility of the solution of these equations leads to specific values of exponential coefficients k in the following form:

$$k_1 = \frac{(-4\omega_0^2 \alpha^4 + \varepsilon^2 \alpha^2 \omega_0^4 + 4\varepsilon^2 \omega_0^6)^{1/2}}{2(\alpha^2 + 4\omega_0^2)} - \frac{4\omega_0^2 \alpha}{2(\alpha^2 + 4\omega_0^2)}, \quad (\text{S4})$$

$$k_2 = -\frac{(-4\omega_0^2 \alpha^4 + \varepsilon^2 \alpha^2 \omega_0^4 + 4\varepsilon^2 \omega_0^6)^{1/2}}{2(\alpha^2 + 4\omega_0^2)} - \frac{4\omega_0^2 \alpha}{2(\alpha^2 + 4\omega_0^2)}. \quad (\text{S5})$$

As a result:

$$\begin{aligned} \varphi_{gshe}(\tau) = (\mu_1 e^{k_1\tau} + \mu_2 e^{k_2\tau}) \text{Cos}(\omega_0\tau) \\ + (\eta_1 \mu_1 e^{k_1\tau} \\ + \eta_2 \mu_2 e^{k_2\tau}) \text{Sin}(\omega_0\tau) \end{aligned} \quad (\text{S6})$$

where

$$\nu_i = -\frac{\alpha k_i + \omega_0^2 \frac{\varepsilon}{2} \text{Cos}(\Delta)}{2\omega_0 k_i + \alpha\omega_0 - \omega_0^2 \frac{\varepsilon}{2} \text{Sin}(\Delta)} \mu_i = \eta_i \mu_i, \quad i = 1, 2. \quad (\text{S7})$$

We assume the particular solution of the inhomogeneous equation (S1) as $\varphi_{psie}(\tau) = A \text{Cos}(\omega_0\tau) + B \text{Sin}(\omega_0\tau)$. The substitution of $\varphi_{psie}(\tau)$ into the equation (S1) leads to the following coupled linear equations for the coefficients A и B :

$$\omega_0 \alpha B + \frac{\omega_0^2 \varepsilon}{2} A \text{Cos}(\Delta) - \frac{\omega_0^2 \varepsilon}{2} B \text{Sin}(\Delta) = i_s, \quad (\text{S8a})$$

$$-\omega_0 \alpha A - \frac{\omega_0^2 \varepsilon}{2} B \text{Cos}(\Delta) - \frac{\omega_0^2 \varepsilon}{2} A \text{Sin}(\Delta) = 0. \quad (\text{S8b})$$

We thus have:

$$A = \frac{2 i_s \varepsilon \text{Cos}(\Delta)}{\omega_0^2 \varepsilon^2 - 4\alpha^2}, \quad (\text{S9})$$

$$B = -\frac{4 i_s \alpha}{(\omega_0^2 \varepsilon^2 - 4\alpha^2) \omega_0} - \frac{2 \varepsilon i_s \text{Sin}(\Delta)}{(\omega_0^2 \varepsilon^2 - 4\alpha^2)}. \quad (\text{S10})$$

Next, we apply the initial conditions in order to determine the coefficients μ_1 and μ_2 :

$$\mu_1 + \mu_2 + A = i_s, \quad (\text{S11a})$$

$$\omega_0(\eta_1 \mu_1 + \eta_2 \mu_2) + B\omega_0 + k_1 \mu_1 + k_2 \mu_2 = 0. \quad (\text{S11b})$$

Solving equations (S11), we find μ_1 and μ_2 to be:

$$\mu_1 = -\frac{(i_s - A)(\omega_0 \eta_2 + k_2) + B\omega_0}{k_1 - k_2 + \eta_1 \omega_0 - \eta_2 \omega_0}, \quad (\text{S12})$$

$$\mu_2 = -\frac{(A - i_s)(\omega_0 \eta_1 + k_1) - B\omega_0}{k_1 - k_2 + \eta_1 \omega_0 - \eta_2 \omega_0}. \quad (\text{S13})$$

Thus, the total solution of the equation (S1) takes the following form:

$$\begin{aligned}\varphi(\tau) = & \text{Cos}(\omega_0\tau)(A + \mu_1 e^{k_1\tau} + \mu_2 e^{k_2\tau}) \\ & + \text{Sin}(\omega_0\tau)(B + \eta_1\mu_1 e^{k_1\tau} \\ & + \eta_2\mu_2 e^{k_2\tau}).\end{aligned}\quad (\text{S14})$$

A non-resonant case.

Similarly, to the approach given in the resonant case, we obtain an analytical solution of a second order inhomogeneous differential equation in the non-resonant case:

$$\begin{aligned}\ddot{\varphi} + \alpha\dot{\varphi} + [1 + \varepsilon \text{Cos}(2\Omega\tau + \Delta)]\omega_0^2\varphi \\ = i_s \text{Cos}(\Omega\tau).\end{aligned}\quad (\text{S15})$$

The general solution of the equation (S15) is the sum (S2) of the general solution of a homogeneous equation and the particular solution of an inhomogeneous equation. We assume the general solution of the homogeneous equation as $\varphi_{gshc}(\tau) = \delta e^{k\tau} \text{Cos}(\Omega\tau) + \rho e^{k\tau} \text{Sin}(\Omega\tau)$. The substitution of $\varphi_{gshc}(\tau)$ into the equation (S15) leads to the following coupled linear equations for the coefficients δ и ρ :

$$\begin{aligned}\delta \left(\alpha\gamma + \omega_0^2 - \Omega^2 + \frac{\varepsilon\omega_0^2}{2} \text{Cos}(\Delta) \right) \\ + \beta \left(2\Omega\gamma + \Omega\alpha - \frac{\varepsilon\omega_0^2}{2} \text{Sin}(\Delta) \right) \\ = 0,\end{aligned}\quad (\text{S16a})$$

$$\begin{aligned}\beta \left(-2\Omega\gamma + \omega_0^2 - \Omega^2 - \frac{\varepsilon\omega_0^2}{2} \text{Cos}(\Delta) \right) \\ + \delta \left(\alpha\gamma - \Omega\alpha - \frac{\varepsilon\omega_0^2}{2} \text{Sin}(\Delta) \right) \\ = 0.\end{aligned}\quad (\text{S16b})$$

The compatibility of the solution of these equations leads to specific values of exponential coefficients γ in the following form:

$$\begin{aligned}\gamma_1 = \frac{1}{\alpha^2 + 4\Omega^2} \left(\alpha^2(\omega_0^2 + \Omega^2)^2 - (\alpha^2 + 4\Omega^2)(\alpha^2\Omega^2 \right. \\ \left. + \Omega^4 - 2\omega_0^2\Omega^2 + \omega_0^4 - \left(\frac{\varepsilon\omega_0^2}{2}\right)^2) \right)^{1/2} \\ - \frac{\alpha(\omega_0^2 + \Omega^2)}{\alpha^2 + 4\Omega^2},\end{aligned}\quad (\text{S17})$$

$$\begin{aligned}\gamma_2 = -\frac{1}{\alpha^2 + 4\Omega^2} \left(\alpha^2(\omega_0^2 + \Omega^2)^2 - (\alpha^2 + 4\Omega^2)(\alpha^2\Omega^2 \right. \\ \left. + \Omega^4 - 2\omega_0^2\Omega^2 + \omega_0^4 - \left(\frac{\varepsilon\omega_0^2}{2}\right)^2) \right)^{1/2} \\ - \frac{\alpha(\omega_0^2 + \Omega^2)}{\alpha^2 + 4\Omega^2}.\end{aligned}\quad (\text{S18})$$

In that way:

$$\begin{aligned}\varphi_{\text{опов}}[\tau] = & (\delta_1 e^{\gamma_1\tau} + \delta_2 e^{\gamma_2\tau}) \text{Cos}(\Omega\tau) \\ & + (\chi_1 \delta_1 e^{\gamma_1\tau} \\ & + \chi_2 \delta_2 e^{\gamma_2\tau}) \text{Sin}(\Omega\tau)\end{aligned}\quad (\text{S19})$$

where

$$\begin{aligned}\rho_i = -\frac{\gamma_i\alpha + \omega_0^2 - \Omega^2 + \frac{\varepsilon\omega_0^2}{2} \text{Cos}(\Delta)}{2\Omega\gamma_i + \alpha\Omega - \frac{\varepsilon\omega_0^2}{2} \text{Sin}(\Delta)} \chi_i = \chi_i \delta_i, \\ i = 1, 2.\end{aligned}\quad (\text{S20})$$

We assume the particular solution of the inhomogeneous equation (S1) as $\varphi_{psie}(\tau) = A_1 e^{i\Omega\tau} + B_1 e^{-i\Omega\tau}$. After the substituting of $\varphi_{gshc}(\tau)$ into the equation (S15) the following coupled linear equations for the coefficients A_1 and B_1 were obtained:

$$\begin{aligned}A_1(\omega_0^2 - \Omega^2 + i\alpha\Omega) + \frac{\varepsilon\omega_0^2}{2} (\text{Cos}(\Delta) + i\text{Sin}(\Delta))B_1 \\ = \frac{i_s}{2},\end{aligned}\quad (\text{S21a})$$

$$\begin{aligned}B_1(\omega_0^2 - \Omega^2 - i\alpha\Omega) + \frac{\varepsilon\omega_0^2}{2} (\text{Cos}(\Delta) - i\text{Sin}(\Delta))A_1 \\ = \frac{i_s}{2}.\end{aligned}\quad (\text{S21b})$$

We thus have

$$\begin{aligned}A_1 = B_1^+ \\ = i_s \frac{\omega_0^2 \varepsilon \text{Cos}(\Delta) - 2(\omega_0^2 - \Omega^2)}{(\omega_0^2 \varepsilon)^2 - 4(\omega_0^2 - \Omega^2)^2 + (\alpha\Omega)^2} \\ + i_s \frac{2\alpha\Omega + \omega_0^2 \varepsilon \text{Sin}(\Delta)}{(\omega_0^2 \varepsilon)^2 - 4((\omega_0^2 - \Omega^2)^2 + (\alpha\Omega)^2)} i \\ = \tilde{A}_1 + \tilde{B}_1 i.\end{aligned}\quad (\text{S22})$$

It is more convenient to introduce new variables defined by $A_0 = 2\tilde{A}_1$ and $B_0 = -2\tilde{B}_1$. In that case:

$$\varphi_{psie}(\tau) = A_0 \text{Cos}(\Omega\tau) + B_0 \text{Sin}(\Omega\tau).\quad (\text{S23})$$

Next, we use the initial conditions in order to determine the coefficients δ_1 and δ_2 :

$$\delta_1 + \delta_2 + A_0 = i_s, \quad (\text{S24a})$$

$$\Omega(\chi_1\delta_1 + \chi_2\delta_2) + B_0\Omega + \gamma_1\delta_1 + \gamma_2\delta_2 = 0. \quad (\text{S24b})$$

from where follows:

$$\delta_1 = -\frac{(i_s - A_0)(\omega_0\chi_2 + \gamma_2) - B_0\Omega}{\gamma_1 - \gamma_2 + \chi_1\omega_0 - \chi_2\omega_0}, \quad (\text{S25})$$

$$\delta_2 = \frac{(i_s - A_0)(\omega_0\chi_1 + \gamma_1) - B_0\Omega}{\gamma_1 - \gamma_2 + \chi_1\omega_0 - \chi_2\omega_0}. \quad (\text{S26})$$

Thus, we obtained an analytical solution of the equation (S15):

$$\begin{aligned} \varphi(\tau) = & (\delta_1\text{Exp}(\gamma_1\tau) + \delta_2\text{Exp}(\gamma_2\tau) + A_0)\text{Cos}(\Omega\tau) \\ & + (\delta_1\chi_1\text{Exp}(\gamma_1\tau) + \delta_2\chi_2\text{Exp}(\gamma_2\tau) \\ & + B_0)\text{Sin}(\Omega\tau). \end{aligned} \quad (\text{S27})$$