

“SUPPLEMENTAL MATERIAL to”

Study of cesium atomic transitions in strong magnetic fields with use of a half-wavelength-thick cell

A. Sargsyan, G. Hakhumyan, R. Mirzoyan, D. Sarkisyan

Hyperfine spectra of Cs D₂ line transitions for zero magnetic field consist of two groups shown in Fig.1: group of atomic transitions $F_g=3 \rightarrow F_e=2,3,4$ (with frequency separation 151 MHz and 201.5 MHz) and group of atomic transitions $F_g=4 \rightarrow F_e=3,4,5$ with frequency separation 201.5 MHz and 251 MHz). The Doppler broadening of each atomic transition is ≈ 400 MHz, that's why inside absorption spectrum for a usual 0.1 - 3 cm-long Cs cell (at a room temperature) these transitions are not resolved. Use of the nanometric- thin cell (NTC), with $L=\lambda/2$ thickness, due to the effect of the spectral narrowing allows one to resolve 6 hyperfine atomic transitions. The experimental profile of the absorption line is best approximated by the curve which is described by the “pseudo-Voigt” function (the “Origin-8” program).

When a strong magnetic field is applied to Cs NTC, there is splitting and shift of Cs D₂ line energy levels. That's why new atomic transitions appeared. In Fig.2 absorption spectrum of the Cs NTC with $L=\lambda/2$, for circular σ^- laser excitation and $B=6.1$ kG is shown. In the case of hyperfine Paschen-Back (HPB) regime there are two groups and each group contains 8 atomic transitions: 1-st group transitions labeled 1'-8' belongs to $6S_{1/2}, m_J=-1/2 \rightarrow 6P_{3/2}, m_J=-3/2$ transitions, while 2-nd group transitions 9'-16' belongs to $6S_{1/2}, m_J=+1/2 \rightarrow 6P_{3/2}, m_J=-1/2$ (see Fig.3). The frequencies of all 16 transitions are below the frequency of $F_g=4 \rightarrow F_e=3$ transition.

It is well known that a radiation with linear polarization in the case of a longitudinal magnetic field could be considered as a sum of two radiations with circular σ^+ and σ^- polarizations. In Fig.4 absorption spectrum of the Cs NTC with $L=\lambda/2$, for linear polarization excitation and $B=6.1$ kG is shown. There are 16 atomic transitions both for σ^+ and σ^- excitations. It is interesting to note that from Fig.4 it is seen that the ratio of the amplitudes for atomic transitions in the case of σ^+ excitation $A(9-16) / A(1-8) \approx 3$, meanwhile in the case of σ^- excitation $A(1'-8') / A(9'-16') \approx 3$ [for the groups 1-8 and 1'-8' the initial level is ($F_g=3$) and for the groups 9-16 and 9'-16' the initial level is ($F_g=4$)]. This is confirmed also by the theory.

Thus, $\lambda/2$ -method is a convenient tool for quantitative study of atomic transitions in magnetic fields.

Figure Captions

Fig.1 Absorption spectra of Cs D_2 line obtained with NTC of $L=\lambda/2$. They consist of two groups: atomic transitions $Fg=3 \rightarrow Fe=2,3,4$ and atomic transitions $Fg=4 \rightarrow Fe=3,4,5$. All six transitions are partially resolved, meanwhile in a usual Cs cell of cm-size they are not resolved at all.

Fig.2 Absorption spectrum of the Cs NTC with $L=\lambda/2$, $B=6.1$ kG, for circular σ^- laser excitation. All sixteen atomic transitions are well resolved and are located below $Fg=4 \rightarrow Fe=3$ transition.

Fig.3. Diagram of the Cs D_2 line atomic transitions for σ^- laser excitation for the case of HPB regime. The selection rules for the atomic transitions are $\Delta m_j = -1$; $\Delta m_l = 0$. It should be 16 atomic transitions. Ground levels ($Fg=3$) and ($Fg=4$) for $B \gg B_0$ contain by 8 sublevels.

Fig.4 Absorption spectrum of Cs NTC with $L=\lambda/2$, $B=6.1$ kG, for linear laser excitation. All 32 atomic transitions are well resolved .16 transitions (for σ^+ excitation) are located above $Fg=3 \rightarrow Fe=4$,while another 16 transitions (for σ^- excitation) are located below $Fg=4 \rightarrow Fe=3$ transition.

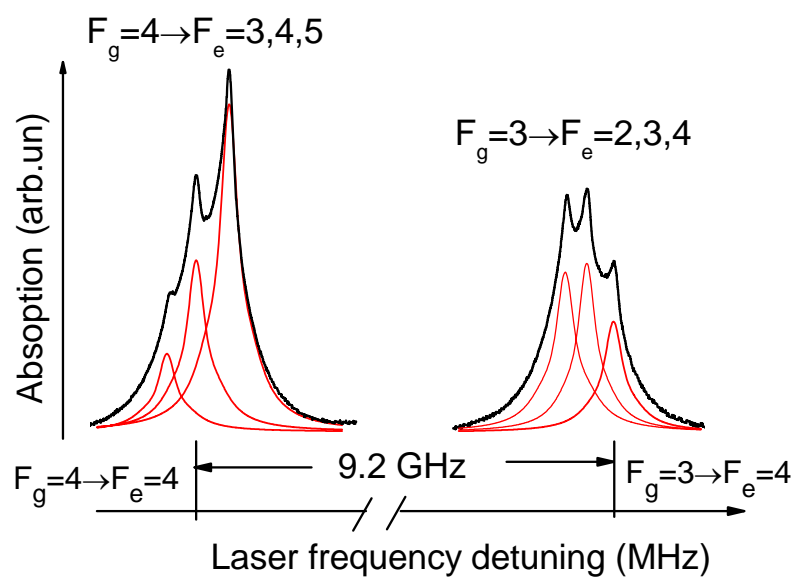


Fig.1 Supplementary

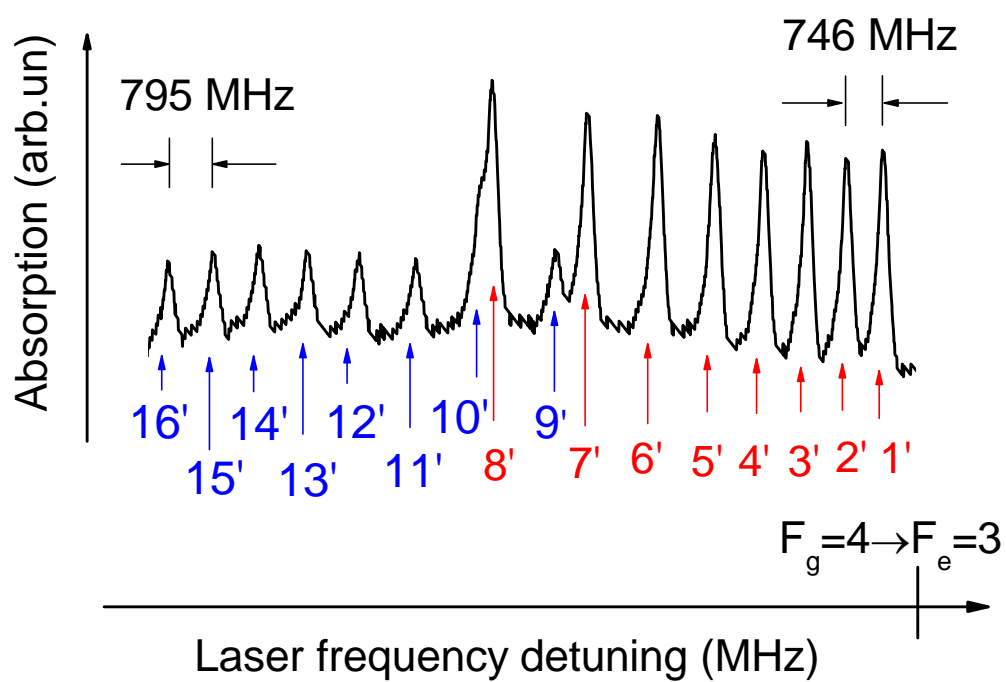


Fig.2 Supplementary

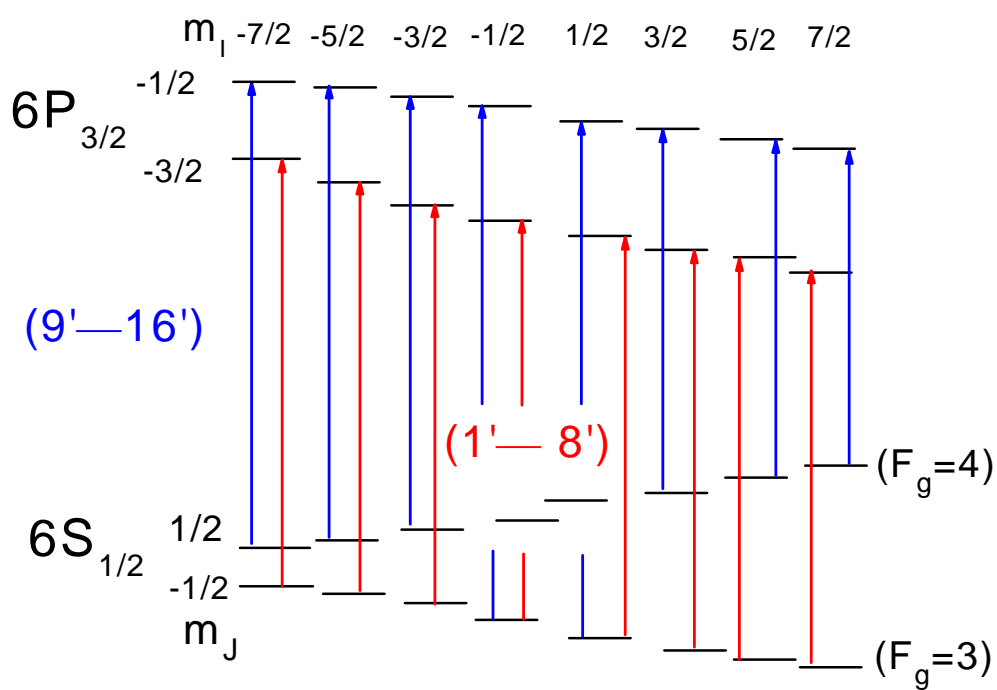


Fig.3 Supplementary

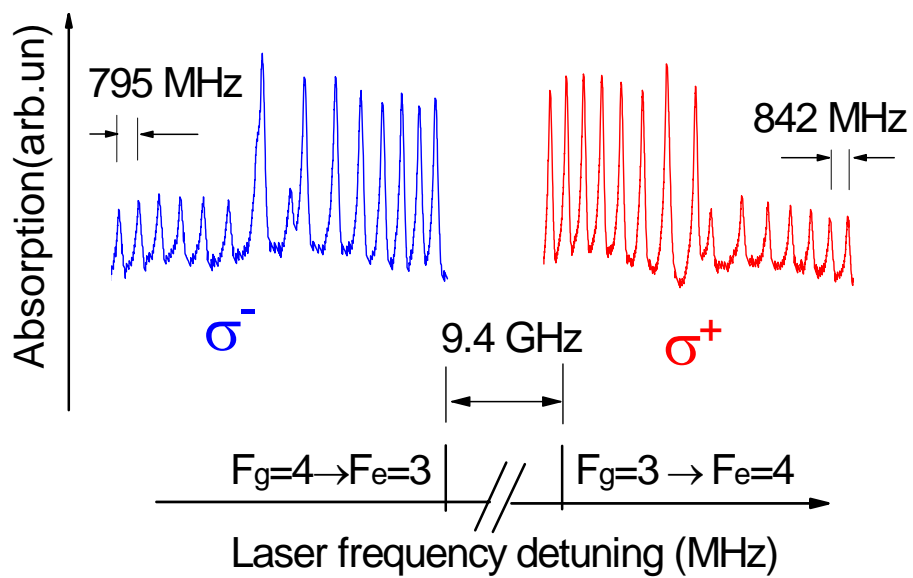


Fig.4 Supplementary